

## NOTATION

$f(x, y)$  = general function of dependent variables  $x$  and  $y$   
 $g(x, y, \epsilon)$  = general function of dependent variables  $x, y$   
 and of parameter  $\epsilon$   
 $t$  = independent variable  
 $t_l$  = boundary layer length  
 $w(x, y)$  = general function of dependent variables  $x$  and  $y$   
 $x$  = nonstiff dependent variable  
 $x_0, x_1$  = nonstiff outer variables  
 $X_0, X_1$  = nonstiff inner variables  
 $y$  = stiff dependent variable  
 $y_0, y_1$  = stiff outer variables  
 $Y_0, Y_1$  = stiff inner variables

## Greek Letters

$\epsilon$  = bookkeeping indication of degree of stiffness  
 $\zeta$  = initial condition on nonstiff variable  
 $\eta$  = initial condition on stiff variable  
 $\alpha$  = dimensionless constant  
 $\lambda_i$  =  $i$ th eigenvalue  
 $\rho$  = spectral radius  
 $\tau$  = expanded inner  $t$  scale,  $t/\epsilon$

## Subscripts

$x$  = partial derivative with respect to  $x$   
 $y$  = partial derivative with respect to  $y$   
 $0$  = zeroth-order expansion in  $\epsilon$   
 $1$  = first-order expansion in  $\epsilon$

## LITERATURE CITED

- Aiken, R. C., and L. Lapidus, "An Effective Numerical Integration Method for Typical Stiff Systems," *AIChE J.*, **20**, 368 (1974).  
 Birkhoff, G., "Numerical Solution of Non-linear Differential Equations," *Proc. Symp. of the Math. Research Center*, U.S. Army, University of Wisconsin, Madison (1966).  
 Bodenstein, M., and H. Lutkemeyer, "Die Photochemische Bildung von Bromwasserstoff, und die Bildungsgeschwindigkeit der Brommolekel aus den Atomen," *Z. Phys. Chem.*, **114**, 208 (1925).  
 Edelson, D., and D. L. Allara, "Parameterization of Complex Reaction Systems: Model Fitting vs. Fundamental Kinetics," *AIChE J.*, **19**, 638 (1973).  
 Gear, C. W., "The Automatic Integration of Ordinary Differential Equations," *Comm. ACM*, **14**, 176 (1971).  
 Gelinas, R. J., "Stiff Systems of Kinetic Equations—A Practitioner's View," *J. Comp. Phys.*, **9**, 222 (1972).  
 Heineken, F. G., H. M. Tschuchiy, and R. Aris, "On the Mathematical Status of the Pseudo-State Hypothesis of Biochemical Kinetics," *Math. Biosci.*, **1**, 95 (1967).  
 Herriott, G. E., R. E. Eckert, and L. F. Albright, "Kinetics of Propane Pyrolysis," *AIChE J.*, **18**, 84 (1972).  
 Hirschfelder, J. O., "Pseudo-Stationary State Approximation in Chemical Kinetics," *J. Chem. Phys.*, **26**, 271 (1957).  
 Johnson, H. S., *Gas Phase Reaction Rate Theory*, The Ronald Press, New York (1966).  
 Keneshea, T. J., "A Technique for Solving the General Reaction-Rate Equations in the Atmosphere," U.S. Air Force Cambridge Research Lab., Environmental Research Paper No. 263 (1967).  
 Luss, D., and N. R. Amundson, "Stability of Batch Catalytic Fluidized Beds," *AIChE J.*, **14**, 211 (1968).  
 Ponzo, P. J., and N. Wax, "Relaxation Oscillations, Parasitics, and Singular Perturbations," *IEEE Trans. Circuit Theory*, CT-19, 623 (1972).  
 Ray, W. H., "The Quasi-Steady-State Approximation in Continuous Stirred Tank Reactors," *Can. J. Chem. Eng.*, **47**, 503 (1969).  
 Snow, R. H., "A Chemical Kinetics Program for Homogeneous and Free-Radical Systems of Reactions," *J. Phys. Chem.*, **70**, 2780 (1966).  
 Tikhonov, A. N., "Systems of Differential Equations Containing Small Parameters in the Derivatives," *Math. Sbornik*, **31**, 73 (1952); (Eng. trans. by A. Muzyka, Dept. of Transportation Systems Center, Cambridge, Mass. USA).

Manuscript received April 9 and accepted April 14, 1975.

# Prediction of "Tortuosity Factors" from Pore Structure Data

F. A. L. DULLIEN

Department of Chemical Engineering  
 University of Waterloo, Waterloo, Ontario, Canada

In a recent paper the author has proposed a new network permeability model of porous media (Dullien, 1975) consisting of a set of cubic networks of arbitrary orientation with respect to the macroscopic flow direction. This model has resulted in the following formula to predict permeabilities:

$$k_m = \frac{\epsilon}{96 \sum_i \sum_j V_{ij}} \sum_i \frac{(\sum_j V_{ij}/D_j^2)^2}{\sum_j V_{ij}/D_j^8} \quad (1)$$

where  $\epsilon$  is the porosity corresponding to the pores considered in the calculations and  $V_{ij}$  is the volume of the pores of diameter  $D_j$ , the entry to which is controlled by pore necks of diameter  $D_i$ . The bivariate pore size distribution  $V_{ij}$  can be estimated from the mercury intrusion porosimetry and the photomicrographic pore size distribution curves of the sample. In the same paper the author has shown that a cubic network of capillary tubes has isotropic permeability which is independent of the orientation of the network. The calculations have shown that the network be-

haves in every orientation as if 1/3 of the tubes were conducting, resulting in the factor  $3 \times 32$  in the denominator of Equation (1). In other treatments (Haring and Greenkorn, 1970; Johnson and Stewart, 1965; Wiggs, 1958), it has been shown that for pores of random orientation also a factor of 1/3 results as compared with the factor 1 for the case when all the pores are lined up parallel to the macroscopic flow direction. Often it has been the custom to call the reciprocal of this factor tortuosity  $T$ .

Equation (1) has been applied successfully to highly compacted materials of very different permeabilities (Dullien, 1975).

In the case of pore diffusion and of electric conductivity measurements in porous media, a different kind of tortuosity  $X$  has been defined by the equation (Wyllie and Rose, 1950; Cornell and Katz, 1953)

$$F = \frac{X}{\epsilon} \quad (2)$$

where

$$F \equiv \frac{D_{12}}{D_{12,\text{eff}}} = \frac{\rho_0}{\rho_w} \quad (3)$$

Here  $D_{12}$  is the binary diffusion coefficient in a mixture of substances 1 and 2, and  $D_{12,\text{eff}}$  is the corresponding effective diffusion coefficient measured in the porous solid.  $\rho_0$  is the specific resistance of the porous medium to the flow of the electric current through a conducting fluid filling the entire pore space and  $\rho_w$  is the resistivity of the conducting fluid in the bulk state.  $F$  is often called the *resistivity factor*.

The tortuosity  $X$  defined by Equation (2), unlike the tortuosity  $T$ , depends on the degree of constrictedness of the capillaries. This can be easily demonstrated on the example of a porous medium consisting of identical parallel capillaries with periodic step changes in their diameter. For this kind of porous medium, the tortuosity  $T$  is equal to 1. For  $X$ , however, the following expression is readily obtained by using Equations (2) and (3):

$$X = \frac{(1 + x^2y)(1 + y/x^2)}{(1 + y)^2} \quad (4)$$

where  $x = D_l/D_s$  and  $y = l_l/l_s$ .  $D_l$  and  $D_s$  is the diameter of the large and small capillary segments, respectively, and  $l_l$  and  $l_s$  are the corresponding lengths. The same expression was derived a long time ago by Michaels (1959). It is evident from Equation (4) that as  $x$  increases without limit,  $X$  will also increase indefinitely. This phenomenon can be probably best understood if one imagines that in every continuous capillary path in a porous medium at some point a tiny orifice is installed whose opening is chosen arbitrarily small as compared with the capillary diameter. By this artifice the porosity and the effective path length would be changed only negligibly; nevertheless the value of  $X$  could be made arbitrarily great. For this reason the quantity  $X$  should not be called *tortuosity* at all because it measures also an effect which is not due to the tortuousness of the paths in the medium. The confusion which has been noted in the literature is exemplified by the following typical statement (Schopper, 1966): "That the rock resistivity exceeds the electrolyte resistivity by the factor  $F$  is caused by two effects: first of all, by the reduction of the cross section available for conduction, from bulk cross section  $A$  to total pore cross section  $\epsilon A$ ; secondly by the fact that the true current path length  $L_e$  exceeds the length  $L$  of the sample." It is obvious from this statement that the reduction of cross section available for conduction has been equated with the average reduction of cross section due to the presence of the solid matrix. The fact, however, that the effective cross section available for conduction can be much smaller than  $\epsilon A$  because of constrictions in the capillaries has been overlooked.

In the literature on heterogeneous catalysis (Satterfield, 1970),  $X$  has been split into the product of two factors

$$X = T \cdot S' \quad (5)$$

where  $T$  accounts for increased path length, and  $S'$  was called *shape factor*. The misnomer *tortuosity factor* for  $X$ , however, is still being used here, too. A better term for  $S'$  might be *constriction factor*.

As successful application of the author's network permeability model has shown (Dullien, 1975), the value of the tortuosity  $T$  can be expected to be equal to 3 or less. The same network model, however, can be used also to predict the value of  $X$  from the bivariate pore size distribution data. The derivation parallels that already published for the permeability calculation; therefore, details will not be given here. For each different controlling segment di-

ameter  $D_i$ , a different network  $i$  is used. The pore size distribution is also different in each network  $i$ . It is assumed that the  $i$  networks are independent and, therefore, their conductivities are additive. The orientation of each network with respect to the direction of the macroscopic flux is arbitrary. The total current  $I_i$  in an  $i$  network is written as the sum of the currents in the capillaries pointing in the  $x$ ,  $y$ , and  $z$  coordinate directions

$$I_i = (I_i)_x + (I_i)_y + (I_i)_z \quad (6)$$

After some algebra, one obtains

$$I_i = \frac{3\pi \Delta E}{4\rho_w \sum_j n_{ij} l_{ij}/D_j^2} (\cot^2\alpha + \cot^2\beta + \cot^2\alpha \cot^2\beta) \quad (7)$$

where  $\Delta E$  is the drop in the potential between the two faces of the sample.  $l_{ij}$  is the length and  $n_{ij}$  the number, respectively, of capillary segments associated with the volume  $V_{ij}$ .  $\alpha$  and  $\beta$  are the two angles defining the orientation of the sample with respect to the macroscopic flux. Using Equation (3) and the relations given in Dullien (1975) for the area of the face, the distance between the two faces of the sample in terms of the angles  $\alpha$  and  $\beta$ , and the void fraction  $\epsilon_i$  corresponding to the volume of the tubes in the  $i$  network, the final expression for  $F_i$  is obtained

$$F_i = \frac{3}{\epsilon} \frac{(\sum_j \epsilon_j V_{ij}) (\sum_j V_{ij}/D_j^4)}{(\sum_j V_{ij}/D_j^2)^2} \quad (8)$$

From the assumption of additive conductivities of the  $i$  networks there follows

$$F = \frac{1}{\sum_i (1/F_i)} \quad (9)$$

Thus

$$F = \frac{(3/\epsilon) \sum_i \epsilon_i V_{ij}}{\sum_i \frac{(\sum_j V_{ij}/D_j^2)^2}{\sum_j V_{ij}/D_j^4}} \quad (10)$$

Comparing this with Equation (2) it is evident that

$$X = \frac{(3) \sum_i \epsilon_i V_{ij}}{\sum_i \frac{(\sum_j V_{ij}/D_j^2)^2}{\sum_j V_{ij}/D_j^4}} \quad (11)$$

where 3 is the value of the tortuosity  $T$  and the rest of the expression is the constriction factor  $S'$ .

Equation (11) permits the prediction of  $S'$  from pore structure data. It is important to point out that whereas Equation (1) for the permeability has the property that the terms corresponding to decreasing controlling pore entry diameters  $D_i$  rapidly become insignificant and thus negligible, this is not the case for Equation (11) obtained for  $X$ . When calculating permeabilities one can safely neglect the smallest pores, but in the calculation of  $X$  the entire range of pore sizes should be covered.

At the present there are no experimental data on  $F$  in the case of materials for which the bivariate pore size distribution has been evaluated. Instead, the predictions of Equation (11) have been used to obtain values for the exponent  $m$  in the time-honored relationship due to Archie (1942)

$$F = \epsilon^{-m} \quad (12)$$

for the 14 sandstone samples for which the bivariate distribution has been calculated. The values obtained for  $X$  range from 7 to 60 and the values of  $m$  range from 2.2 to 3.0. In the majority of the cases  $X$  is around 10 and  $m$  is

around 2.5. These values are within the range reported for tightly compacted materials.

## NOTATION

$A$  = bulk cross section of sample  
 $D_i$  = controlling pore entry diameter  
 $D_j$  = pore diameter  
 $D_l$  = diameter of large pore segment  
 $D_s$  = diameter of small pore segment  
 $D_{12}$  = binary diffusion coefficient  
 $D_{12,eff}$  = effective diffusion coefficient  
 $E$  = electric potential  
 $F$  = quantity defined by Equation (3)  
 $I$  = current  
 $k_m$  = permeability predicted by Dullien's model  
 $l_{ij}$  = length of a pore segment associated with the volume  $V_{ij}$   
 $l_l$  = length of a large pore segment  
 $l_s$  = length of a small pore segment  
 $L$  = length of sample  
 $L_e$  = effective path length in sample  
 $m$  = quantity defined in Equation (12)  
 $n_{ij}$  = number of pore segments associated with the volume  $V_{ij}$   
 $S'$  = quantity defined by Equation (5)  
 $T$  = tortuosity  
 $V_{ij}$  = bivariate pore volume distribution  
 $x$  =  $D_l/D_s$   
 $X$  = quantity defined by Equations (2) and (3)  
 $y$  =  $l_l/l_s$

## Greek Letters

$\alpha$  = angle defining orientation of sample with respect to direction of macroscopic current

$\beta$  = angle defining orientation of sample with respect to direction of macroscopic current  
 $\epsilon$  = porosity  
 $\rho_0$  = specific resistivity of porous medium saturated with electrolyte  
 $\rho_w$  = specific resistivity of electrolyte

## LITERATURE CITED

- Archie, G. E., "The Electrical Resistivity Log as an Aid in Determining Some Reservoir Characteristics," *Trans. AIME*, **146**, 54 (1942).  
 Cornell, D., and D. L. Katz, "Flow of Gases through Consolidated Porous Media," *Ind. Eng. Chem.*, **45**, 2145 (1953).  
 Dullien, F. A. L., "New Network Permeability Model of Porous Media," *AIChE J.*, **21**, 299 (1975).  
 Haring, R. E., and R. A. Greenkorn, "A Statistical Model of a Porous Medium with Non-Uniform Pores," *AIChE J.*, **16**, 477 (1970).  
 Johnson, M. F. L., and W. E. Stewart, "Pore Structure and Gaseous Diffusion in Solid Catalysis," *J. Catalysis*, **4**, 248 (1965).  
 Michaels, A. S., "Diffusion in a Pore of Irregular Cross Section—A Simplified Treatment," *AIChE J.*, **5**, 270 (1959).  
 Satterfield, C. N., *Mass Transfer in Heterogeneous Catalysis*, M.I.T. Press, Cambridge, Mass. (1970).  
 Schopper, J. R., "A Theoretical Investigation on the Formation Factor/Permeability/Porosity Relationship Using a Network Model," *Geophys. Prospecting*, **14**, (3), 301 (1966).  
 Wiggs, P. K. C., "The Relation between Gas Permeability and Pore Structure of Solids," in *The Structure and Properties of Porous Materials*, p. 183, D. H. Everett and F. S. Stone (eds.), Academic Press, New York (1958).  
 Wyllie, M. R. J., and W. D. Rose, "Some Theoretical Considerations Related to the Quantitative Evaluation of the Physical Characteristics of Reservoir Rock from Electrical Log Data," *Trans. AIME*, **189**, 105 (1950).

Manuscript received March 4 and accepted April 4, 1975.

# A Solution for the Oxygen Mass Transfer Problem in Immobilized Enzyme Systems

FRANCOIS LAWNY, MICHEL CORDONNIER, and DANIEL THOMAS

Laboratory of Enzyme Technology  
 Université de Technologie de Compiègne, France

The kinetic limitations due to oxygen transfer problems were widely described in chemical and biochemical engineering (Ryu et al., 1973; Topiwala and Hamer, 1973). This subject was especially studied in the field of industrial fermentors. In the growing field of immobilized enzymes, not only has the problem not been solved, but also it has never been clearly studied. Numerous papers (Weetall, 1974) have been devoted to immobilized enzymes and recently reviewed (Gryszkiewicz, 1971; Goldman, 1971).

The aim of this note is to show that in immobilized enzyme systems the oxygen can be generated inside the support itself through another enzyme reaction. The basic principle of the described solution is the existence of substrate and product concentration profiles inside the support (Thomas et al., 1974; Barbotin and Thomas,

1974). In the experimental system glucose-oxidase (E.C. 1.1.3.4.—Sigma) and catalase (E.C. 1.1.1.1.—Sigma) are immobilized inside proteic membranes with a co-cross-linking method previously described (Broun et al., 1973). The glucose-oxidase catalyses the transformation of glucose into gluconic acid and  $H_2O_2$  in presence of oxygen. Its kinetics were monitored at pH 8, the optimum pH for the immobilized enzyme, in a phosphate solution  $5.10^{-3}M$  by injection of NaOH in with a pH-stat (Metrohm). Due to the pH value and the presence of lactonase, the lactone decomposition is not the limiting stop of the overall reaction. Under steady state conditions, the enzyme activity is given by the measurement of the  $H^+$  production.

The catalase transforms  $H_2O_2$  into water and oxygen. The reaction rate was followed spectrophotometrically at 240 nm. In both experiments performed with free and immobilized catalase and glucose oxidase, activities were respectively 7500 I.U. (with  $H_2O_2$   $10^{-2}M$ ) and 250 I.U. (with a glucose solution of 20 g/l) 100-ml bath solution. The Thiele modulus values for the catalase are 186 and 5.8 for the glucose-oxidase. The diffusion coefficients were

Correspondence concerning this note should be addressed to D. Thomas. F. Lawny is at the Institut de Technologie des Surfaces and M. Cordonnier is with the Division Fermentation, B.P. 233 Université de Technologie de Compiègne, Compiègne, France.